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## LETTER TO THE EDITOR

## A generalised self-avoiding walk

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#### Abstract

We study a generalised self-avoiding walk on a lattice in which each vertex may be visited less than $k$ times. Turban has argued that, for this model, the critical exponents should change continuously with $k$ from the standard self-avoiding walk exponents ( $k=2$ ) to mean-field exponents. By generating series expansions for both two- and threedimensional lattices, we find strong support for the conventional, and opposing, view that the critical exponents remain unchanged for finite $k$, only taking on their mean-field values in the limit $k \rightarrow \infty$.


Recently Turban (1983) suggested a generalised self-avoiding lattice walk in which the usual self-avoiding constraint does not apply until a particular vertex has been visited $k-1$ times. A $k$ th visit is forbidden. For $k=2$ the usual self-avoiding walk (saw) is recovered, while in the limit $k \rightarrow \infty$ it is clear that the pure random walk is obtained. Intermediate values of $k$ therefore interpolate between these two limits.

During 1983 we investigated this problem, by series analysis, but did not publish our results as they were indicative of the behaviour we expected, which was that the critical exponents do not change with $k$ (for finite $k$ ), and so the generalised problem remains in the same universality class as the saw problem, a result we considered unsurprising.

Turban (1983), however, has given arguments to suggest totally different behaviour. His arguments, based both on a Flory theory argument and on the properties of fractals, suggest that the critical exponent characterising the mean square end-to-end distance through the relation $\left\langle R_{N}^{2}\right\rangle \sim a N^{2 \nu}$, should vary with $k$ like

$$
\begin{equation*}
\nu_{k}=(k+1) /[(k-1) d+2] \tag{1}
\end{equation*}
$$

whenever the spatial dimension $d \leqslant d_{\mathrm{c}}(k)=2 k /(k-1)$.
This surprising result has prompted us to extend our series data and analyse it more carefully. Our re-analysis confirms our earlier results, which are therefore totally different from those obtained by Turban.

The problem appears to have been first suggested by Malakis (1975), who defined the model in passing and suggested the results found in this paper, that is, that there is no change in the critical exponents from the saw values. In a subsequent paper, Malakis (1976) called these walks ' $k$-tolerant walks', where ' $k$ ' in Malakis's notation is ' $k-1$ ' in the notation of Turban, and this paper. Malakis suggested that the problem is essentially different to the saw problem in three dimensions (but not one and two) by invoking an argument based on the properties of pure random walks, and hence
raised the possibility that the exponents may differ in the three-dimensional case. We shall show that this does not appear to be the case.

Firstly, let us look in more detail at the predictions of Turban's results. For $d=2$, $d<d_{\mathrm{c}}(k)$ for all finite $k$, and $\nu_{k}=(k+1) / 2 k$, so that $\nu_{2}=\frac{3}{4}, \nu_{3}=\frac{2}{3}, \nu_{4}=\frac{5}{8}, \nu_{5}=\frac{3}{5}$, etc. Thus $\nu$ decreases monotonically with increasing $k$ from the SAW value ( $\nu=\frac{3}{4}$ ) to the mean-field value ( $\nu=\frac{1}{2}$ ). Consider now the 'susceptibility' exponent $\gamma$ characterising the growth in the number of chains $C_{n}$ of length $n$ through $C_{n} \sim A \mu^{n} n^{\nu-1}$, where $\mu$ is a latticedependent growth constant. Clearly $\gamma$ must also decrease from its value at $k=2$ ( $\gamma=1.34375$ ) to the mean-field value of $\gamma=1$ as $k \rightarrow \infty$. Analogy with Turban's results for $\nu$ suggest that this decrease should be monotonic in $k$. Our findings are that $\gamma$ remains constant for finite $k$ at the $k=2$ (sAw) value, as might be expected from the following heuristic argument due to C J Thompson (private communication). Drawing on the analogy between the $n>0$ and the $n=0$ realisation of the $n$-vector model, we would expect that setting $k>2$ in the generalised saw model corresponds to increasing the range of the interaction in an $n$-vector model, while still keeping the range finite. Under such circumstances one expects the exponents to remain unchanged.

For the two-dimensional triangular lattice we have generated series expansions for the number of $n$-step chains and for a range of values of $k$. The triangular lattice was chosen in order that the characteristic alternation of loose-packed lattice series would be absent. In table 1 we display the series coefficients for $k=2$ (the saw problem) to $k=6$. Note that, in each case, for $0 \leqslant n<2 k-2$, the coefficients are given by $6^{n}$, which is the pure random walks result. Therefore only for $n \geqslant 2 k-2$ do the series coefficients reflect the self-avoiding constraint, and for that reason, as $k$ increases, an increasing number of series coefficients are needed for the series to be analysed with confidence. Accordingly, we concentrate our analysis on those series with low values of $k$.

For $k=2$ we have the saw series, for which 18 terms are known, and for which we find (Guttmann 1984) $\mu=4.1508$ and $\gamma=1.34375$ (Nienhuis 1982). A simple dlog Padé analysis of the series for $k>2$ (not shown) gives $\mu \approx 5.306, \gamma \approx 1.39$ for $k=3$, $\mu \approx 5.68$ and $\gamma \approx 1.4$ for $k=4$. These results immediately cast doubt on Turban's results, as they superficially suggest that $\gamma$ increases with $k$, rather than decreasing as

Table 1. Chain generating function coefficients for generalised SAw on the triangular lattice.

| $n \backslash k$ | 2 | 3 | 4 | 5 | $\infty$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | 1 | 1 | 1 | 1 | 1 |
| 1 | 6 | 6 | 6 | 6 | 6 |
| 2 | 30 | 36 | 36 | 36 | 36 |
| 3 | 138 | 216 | 216 | 216 | 216 |
| 4 | 618 | 1260 | 1296 | 1296 | 1296 |
| 5 | 2730 | 7206 | 7776 | 7776 | 7776 |
| 6 | 11946 | 40650 | 46440 | 46656 | 46656 |
| 7 | 51882 | 227256 | 276054 | 279936 | 279936 |
| 8 | 224130 | 1262832 | 1633848 | 1678320 | 1679616 |
| 9 | 964134 | 6983730 | 9633366 | 10051782 | 10077696 |
| 10 | 4133166 | 38470220 | 56616140 | 60132100 | 60466176 |
| 11 | 17668938 | 211220800 | 331847200 | 359296300 | 362797056 |
| 12 | 75355206 | 1156490000 | 1940717000 | 2144325000 | 2176782336 |
| 13 | 320734686 | 6317095284 | 11327957196 |  |  |
| 14 | 1362791250 | 34435495872 | 66010769382 |  |  |

Turban's theory would suggest. The reason for this apparent increase seems likely to be due to the presence of confluent singularities, and we investigate this in the following manner. We assume that the general coefficient $C_{n}$ behaves in the following manner for sufficiently large $n$ :

$$
\begin{equation*}
C_{n} \sim A \mu^{n} n^{8}\left(1+B / n^{\Delta}+C / n+D / n^{\Delta+1}+\ldots\right) \tag{2}
\end{equation*}
$$

where $g=\gamma-1$, and $\Delta$ is the exponent corresponding to the confluent singularity. Then the ratio of successive terms is
$r_{n}=C_{n} / C_{n-1} \sim \mu\left[1+g / n-B \Delta / n^{1+\Delta}+O\left(n^{-2}\right)+O\left(n^{-3 \Delta}\right)+O\left(n^{-2-\Delta}\right)+O\left(n^{-1-2 \Delta}\right)\right]$.
To eliminate $\mu$, we form ratios of ratios, and find

$$
\begin{equation*}
s_{n}=r_{n} / r_{n-1}-1 \sim-g / n^{2}+h / n^{2+\Delta}+\mathrm{O}\left(n^{-3}\right)+\mathrm{O}\left(n^{-2-2 \Delta}\right) \tag{4}
\end{equation*}
$$

where $h$ can be expressed in terms of $\Delta, g$ and $B$. If $\gamma$ (and hence $g$ ) is known, the quantity

$$
\begin{equation*}
t_{n}=s_{n}+g / n^{2} \sim h / n^{2+\Delta}+\mathrm{O}\left(n^{-3}\right) \tag{5}
\end{equation*}
$$

A plot of $\ln \left|t_{n}\right|$ against $\ln n$ should then produce a straight line with gradient $\Delta+2$. If the value of $g$ used is inaccurate, say $g^{*} \neq g$, the straight line should then have gradient 2 , with intercept giving $g-g^{*}$. In figure 1 we show the $\log -\log$ plot for $k=2$ and 3 , with $g$ chosen to be 0.34375 , in agreement with the $k=2$ (SAW) result of Nienhuis (1982). The plots display decreasing curvature with increasing $n$, and from the obvious trend we obtain, from the last two points, that $\Delta_{3}>0.11(k=3)$ and $\Delta_{2}>0.26(k=2)$. Further, for $k=3$ we have repeated the analysis using $g^{*}=0.23$, and find that the last three points lie on a straight line of gradient 2.0 , the intercept of which suggests $g-g^{*} \approx 0.16$, in agreement with the Padé approximant analysis cited earlier.


Figure 1. Log-log plot of $t_{n}$ (equation (5)) against $n$ for $k=2(\odot)$ and $k=3(\nabla)$ generalised saws on the triangular lattice. Gradient is a measure of $-(2+\Delta)$.

For $k=4$ our series are too short to be successfully analysed by this technique, though the log-log plot of $t_{n}$ against $n$ does have gradient greater than 2 , implying $\Delta_{4}>0$.

While this analysis is by no means conclusive, we believe that the fact that elementary methods of series analysis suggest $\gamma(k=3,4)>\gamma(k=2)$, in addition to the analysis given here, which gives a consistent interpretation that $\gamma(k<\infty)=\gamma(k=2)$, would seem to rule out $\gamma(k>2)<\gamma(k=2)$, as suggested by Turban's results.

Our results based on three-dimensional data lend stronger support to this view. In three dimensions, Turban predicts that only for $k=2$ is $d<d_{c}$, while for $k=3, d=3$ is the critical dimension. At the critical dimension we would expect mean-field exponents, modified by confluent logarithmic singularities. That is, for $d=3, k=3$ Turban predicts $\nu=\frac{1}{2}$, while for $d=3, k=2$ we expect $\nu \approx 0.59$ from rg results and series analysis. To check this result, we have generated series on the simple cubic

Table 2. Chain generating function coefficients and mean square end-to-end distances for simple cubic lattice $k=3$ sAws.

| $n$ | $C_{n}$ | $\left\langle R_{n}^{2}\right\rangle$ |
| ---: | ---: | ---: |
| 0 | 1 | 1.00000000 |
| 2 | 36 | 2.00000000 |
| 3 | 216 | 3.00000000 |
| 4 | 1260 | 4.11428571 |
| 5 | 7350 | 5.23183673 |
| 6 | 42462 | 6.41288682 |
| 7 | 245664 | 7.58909730 |
| 8 | 1412568 | 8.82093889 |
| 9 | 8131062 | 10.04690679 |
| 10 | 46617096 | 11.31987269 |
| 11 | 267490638 | 12.58654861 |
| 12 | 1530512418 | 13.89444664 |
| 13 | 8763095548 | 15.19588758 |



Figure 2. Log-log plot of $\left\langle R_{N}^{2}\right\rangle_{k=2} /\left\langle R_{N}^{2}\right\rangle_{k=3}$ against $N$ for $k=3$ generalised SAws on the simple cubic lattice. Gradient is a measure of $2\left(\nu_{2}-\nu_{3}\right)$.
lattice for $k=3$. We have obtained both the number of chains and their mean square end-to-end distance. If Turban is correct, the ratio $\left\langle R_{N}^{2}\right\rangle_{k=2} /\left\langle R_{N}^{2}\right\rangle_{k=3} \sim C N^{2\left(\nu_{2}-\nu_{3}\right)} \approx$ $C N^{0.18}$. A $\log -\log$ plot of this ratio should be linear with slope given by $2\left(\nu_{2}-\nu_{3}\right)$. The raw data are shown in table 2, and the log-log plot is shown in figure 2. It can be seen from figure 2 that the gradient is decreasing with increasing $N$, and the last two points already yield a straight line with gradient 0.074 .

Assuming this trend continues then yields $\nu_{2}-\nu_{3}<0.37$, compared with Turban's prediction of 0.09 . The behaviour we observe is entirely consistent with the expected behaviour $\nu_{2}=\nu_{3}$, and the previously expressed view that, for $k=3$, the model has the same critical exponents as for $k=2$. Thus in conclusion we find that the generalised saw belongs to the same universality class as the usual saw, contrary to Turban's arguments.

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